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(Russian Original Vol. 26, 1956, Journal of Technical Physics (USSR)

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No. 1

Page	Line	Reads	Should read
1	6	. . . a layer of sulfurous lead;	. . . a layer of lead sulfide.
3	last line	. . . the measurements were confined to 1.3°K.	. . . the measurements were extended down to 1.3°K.
26	line 2 from bottom	Any considerably stronger diffusion effects must affect the dependence of $\bar{\sigma}$ upon H	Diffusion effects should have a much greater influence on the dependence of $\bar{\sigma}$ on H.
28	line 19	. . . $\frac{1}{l_H}$ is appreciable at $\frac{1}{l_T} = A \frac{dT}{dy}$.	. . . $\frac{1}{l_H}$ is replaced by $\frac{1}{l_T} = A \frac{dT}{dy}$.
57	line 23	Or, if lithium pairs are introduced into an antimony layer. . . .	Or, if an antimony layer is treated with lithium vapor . . . .
57	line 2 from bottom	. . . the elasticity of lithium vapor . . .	. . . the vapor pressure of lithium . . .
202	34, 35	. . . diodes . . .	. . . dynodes . . .

No. 2

287	line 7, onward	. . . non-basic carriers . . .	. . . minority carriers . . .
	21	. . . Ziener effect . . .	. . . Zener effect . . .
332	14	. . . admitted the error, considering them lowered by one.	. . . made the mistake of supposing they were reduced by unity.
341	[3]	J. J. Tomson . . .	J. J. Thomson . . .
	[4]	L. Leb . . .	L. Loeb . . .
371	[6]	J. Barden . . .	J. Bardeen . . .
	[7]	Sterngeimer . . .	Sternheimer . . .
378	[2]	St. Meuer . . .	St. Meyer . . .
381	17	. . . special points . . .	. . . singular points . . .
385	line 9	. . . the diameters of the comb structure . . .	. . . the dimensions of the comb structure . . .

Page	Line	Reads	Should read
391	line 13	. . . Babinier . . .	. . . Babinet . . .
	line 22	. . . ($l_{\text{lim}} \ll L$) . . .	. . . ($l_{\text{boundary}} \ll L$) . . .
402	line 5	. . . but because this integral is a function of $\underline{s}$ . . .	. . . but, as a function of $\underline{s}$, this integral . . .
407	line 2 from bottom	. . . inverse lattice node (this intensity is determined by the vector **k** . . .	. . . reciprocal lattice determined by the vector **k** . . .
No. 3			
514	lines 1, 4	. . . Matier . . .	. . . Mathieu . . .
514	line 3	. . . intrinsic values . . .	. . . eigenvalues . . .
516	[4]	MacLaughlan, Theory and Application of Matier Functions	McLachlan, Theory and Applications of Mathieu Functions
562	line 13	. . . a singular relationship . . .	. . . a single-valued relation . . .
No. 4			
745	last line and top of 748	Although it might be expected that the irregular surface of the twisted conductor should have a strong effect on the corona starting voltage, in point of fact it is found that both characteristics are quite similar; calculations of the appropriate electric fields [3] show that it is the weak-field region which determines the value of the corona current.	The agreement between the characteristics is due to the fact that the twisted surface of the conductor has a strong effect on the corona starting voltage, but, as has been shown by calculations of the corresponding electrostatic field [5], has no effect on the weak field which determines the corona current.
748	lines 16 and 17	Hence the only difference in the voltage-current characteristics of split and solid conductors is the difference in the corona starting voltage.	Hence the only difference in the voltage-current characteristics of split and solid conductors is the difference in the corona starting voltage (for small distances between components of the split conductor).
811	last 2 lines	It should be noted that in the case when $n_0 \gg 1$ and the equation $\mathrm{Th}^2 \frac{\pi}{2} \eta_0 \simeq 1 - 4e^{-\pi \eta_0}$, can be simplified and acquires the form	It should also be noted, that in the case where $n_0 \gg 1$ and the approximate equation $\tanh^2 \pi/2\, \eta_0 \simeq \simeq 1 - 4e^{-\pi \eta_0}$ can be used, the formula (12) is simplified and acquires the nature
813	line 11	. . . Ritts method . . .	. . . Ritz method . . .
817	[3]	E. Ianke	E. Jahnke . . .
No. 5			
968	[4]	. . Rossi and Shtaub . . .	. . . Rossi and Staub . . .
993	line 8	. . . for a column which is finite in diameter . . .	. . .for the final diameter of the column . . .
	line 13	. . . in a column with finite diameter there is found an exponential dependence . . .	. . . the final diameter of the column has a power-law dependence . . .

Page	Line	Reads	Should read
994	line 13	In the second case, with a plate having a hole diameter 0.5 mm, the current is some 6% greater.	The current in the latter case is 6% greater than with a plate of hole diameter 0.5 mm.
998	line 26	The nonequilibrium is conditioned by the final speed of loading.	The nonequilibrium is due to the finite speed of loading.
1000	line 2	... setting down the integral function...	. . . representing the integrand . . .
1001	line 4	. . . the appearance of the stress tensor will also depend on the relaxation processes.	. . . the form of the stress tensor will also affect the relaxation processes.
	line 21	. . . where σ^* is the stress over the elapsed time for the relaxation.	. . . where σ^* is the stress at the end of one relaxation period.
1003	line 18	In this case the nonintrinsic integral obtained after integration of Eq. (2) comes down to . . .	In this case the improper integral obtained from integration of Eq. (2) reduces to . . .
1007	lines 1, 2	. . . ϵ . . .	. . . $\dot{\epsilon}$ . . .
	line 21	. . . τ_0 . . .	. . . τ_2 . . .
No. 6			
1133	[18]	K. Herring . . .	C. Herring . . .
1164	[13]	Karslaiv and Eger . . .	Carslaw and Jaeger . . .
No. 7			
1384	line 11	. . . due to the "stretching" of the diffusion lines of nonbasic carriers [4].	. . . because of the stretching of the diffusion lengths for minority carriers [4].
1385	line 9	. . . the Green function Eq. (6) which is space dependent for the limiting conditions (7).	. . . the time-dependent Green's function of Eq. (6) for the boundary conditions (7).
	line 11	. . . single instantaneous source . . .	. . . unit instantaneous source. . .
1387	line 6	. . . in (16) dominate the first term . . .	. . . the first term in (16) is dominant . . .
	line 10	. . . the comparable magnitude has a large number of terms.	. . . a large number of terms are of comparable magnitude.
	[5]	Kh. S. Carslow, Theory of Thermal Conductivity. . .	H. S. Carslow, Theory of Heat Conduction. . .
1420	line 23	. . . reverse isometry . . .	. . . rotational isomerism . . .
	line 26	. . . reverse isomers . . .	. . . rotational isomers . . .
No. 8			
1606	last 2 lines	Here q_α comprises $\mathbf{q}$ . . .	Here q_α is the component of $\mathbf{q}$ . . .

Page	Line	Reads	Should read
1611	line 3	pure liquids	true liquids
	[3]	G. Bete and A. Zommerfeld . . .	H. Bethe and A. Sommerfeld . . .
1758	[3]	D. Garrison, R. Lord and D. Lufburov	D. Harrison, R. Lord and D. Loofbourow
1759	[26] and [33]	G. Gertsberg	G. Hertzberg
	[38]	R. V. Pol', Introduction of Optic ...	R. V. Pohl, Introduction to Optics. . .
1791	line 22 and 23	. . . made it necessary for us to calculate its approximate variation theoretically.	. . . made it necessary for us to use the approximate dependence given by the theory of corresponding states.
No. 9			
1845	[12]	U. Kedi . . .	W. Cady. . .
1901	Eq. (2.5)	a =	a' =
1906	last equation	$\vartheta_0 = \vartheta_{m1} + \vartheta_{m1} + \vartheta_{m2} + \vartheta_{m1}$	$\vartheta_0 = \vartheta_{m1} + \vartheta_{s1} + \vartheta_{m2} + \vartheta_{s2}$
1910	[16]	D. Stoker . . .	J. Stoker . . .
1925	[5]	L. M. Mili-Tomson . . .	L. M. Milne-Thomson . . .
1941	[3]	D. Kai and T. Labi . . .	D. Kaye and T. Laby. . .
1969	[11]	K. F. Gertsfel'd . . .	K. F. Herzfeld . . .
	[16]	F. Zeits. . .	F. Seitz . . .
1993	[3]	V. L. Diurend . . .	V. L. Durand . . .
1995	line 21	. . . during the course of the time . . .	. . . after the time . . .
2046	line 2	. . . decomposition of ternary alloys to binary alloys . . .	. . . decomposition of solid solutions of binary alloys . . .
2047	line 6	. . . is defined by the ratio of coefficients of the p's in (2b) . . .	. . . is the determinant of the coefficient of the p's in (2b).
No. 10			
2101	[7]	Deshman . . .	Dushman. . .
2113	[8]	F. Hand . . .	F. Hund . . .
2192	line 2	. . . the possible meaning of $\underline{n}$ . . .	. . . the possible values of $\underline{n}$ . . .
2195	line 6	. . . with zero initial velocity, its trajectory . . .	. . . with zero initial velocity, and since the limiting curve is an equipotential, its trajectory . . .
2198	line 11	The electrons of . . .	The conversion electrons of . . .

Page	Line	Reads	Should read
2216	[2]	A. Rusterkhol'ts . . .	A. Rusterholz . . .
2298	Table heading, line 2	Average over two hours of operation.	Averaged over hours of operation.
No. 11			
2334	line 6 from bottom	Puncture occurs when $M - \infty$ . . .	[Puncture occurs when $M \rightarrow \infty$ . . . . .
2339	line 6 from bottom	. . . further, the electron concentration at $n = 7 \cdot 10^{17}$ did not depend on the quantity present.	. . . beyond this point, the electron concentration remains equal to $7 \cdot 10^{17}$, independent of the amount of lead excess.
2340	line 7	. . . and as this falls . . .	. . . and then decreases it . . .
2343	line 29	. . . were completely reproducible . . .	. . . were completely reversible . . .
2407	line 26	. . . with colloidal centers which are colored in an electric field the crystals were cooled in air.	. . . with colloidal centers, crystals colored in an electric field were cooled in air.
2455	[10]	O. D. Khvol'son . . .	O. D. Chwolson . . .
No. 12			
2557	footnote	In making approximate calculations we may substitute infinite limits in the integration over ρ in (8) for the limits Ei/p for integration over E/p (10).	In making approximate calculations we may replace the infinite limit in the ρ integration of (8) by the limit E_i/p in the integration over E/P in (10).
2558	line 7	. . . which is almost three times more than . . .	. . . which is almost one-third . . .
2565	[7]	C. Kent. . .	C. Kenty. . .
	[8]	. . . R. Zurman . . .	. . . S. Schurman . . .
2675	line 9	. . . a discontinuity . . .	. . . continuity . . .
	line 29	. . . that only one equation of (3) is sufficient . . .	. . . that Eq. (3) alone is insufficient . . .
	line 31	. . . are determined uniquely.	. . . are not determined uniquely.
2676	line 1	This statement becomes obvious if it is expressed by saying that the bounding surface Σ serves as an ideally conducting boundary of an endovibrator, excited from within by the surface magnetic current i_m.	This statement becomes obvious if we consider that the closed surface Σ acts as an ideally conducting boundary of a cavity oscillator.
2692	line 15	. . . photons . . .	. . . phonons . . .

(See also Errata in No. 5, page 1109)

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(Index of the Journal of Technical Physics of the USSR,

Vol. 26, 1956)

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